Erratum: Goos-Hänchen shift in negatively refractive media [Physical Review E 66, 067603 (2002)]

P. R. Berman (Received 5 January 2005; published 16 March 2005)

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There were some typographical errors in this paper. The lateral displacement *d* is given by $d = -\sin \theta \partial \phi_R / \partial k_x = (1/k_1) \partial \phi_R / \partial \theta$, where $\phi_R = \tan^{-1}[2\alpha \cos \theta / (\cos^2 \theta - \alpha^2)]$. A factor $[1 - (\mu_2 \epsilon_2 / \mu_1 \epsilon_1)]$ was inadvertantly left out of the equation for *d*. The correct expression is

$$d = \frac{2}{k_1} \frac{\left(1 - \frac{\mu_2 \epsilon_2}{\mu_1 \epsilon_1}\right) |\mu_2/\mu_1| \sin \theta}{\left[(\mu_2/\mu_1)^2 \cos^2 \theta + \sin^2 \theta - \frac{\mu_2 \epsilon_2}{\mu_1 \epsilon_1}\right] \sqrt{\sin^2 \theta - \frac{\mu_2 \epsilon_2}{\mu_1 \epsilon_1}}$$

Although the expression for *d* is finite at $\theta = \pi/2$, the derivative $d\phi_R/d\theta$ is not defined at that point, since ϕ_R varies linearly for $\theta < \pi/2$, but is zero for $\theta > \pi/2$. Thus this expression cannot be used for $\theta = \pi/2$, where *d* should equal zero. Renard's [1] expression for *d* is multiplied by $\cos^2 \theta/[1-(\mu_2\epsilon_2/\mu_1\epsilon_1)]$ and vanishes at $\theta = \pi/2$, a point he used to justify his approach, but the approximations he made to arrive at his result were questioned by Lai *et al.* [2].

I would like to thank Dr. S. R. Seshadri for pointing out the missing factor in the expression for *d*. He also noted that Felson [3] obtained "backward" refraction for lateral waves in a calculation of emission at an anisotropic plasma interface.

- [2] H. M. Lai, F. C. Cheng, and W. K. Tang, J. Opt. Soc. Am. A 3, 550 (1986). (The page number has been corrected.)
- [3] L. B. Felson, IEEE Trans. Antennas Propag. 10, 347 (1962).

^[1] R. H. Renard, J. Opt. Soc. Am. 54, 1190 (1964).